

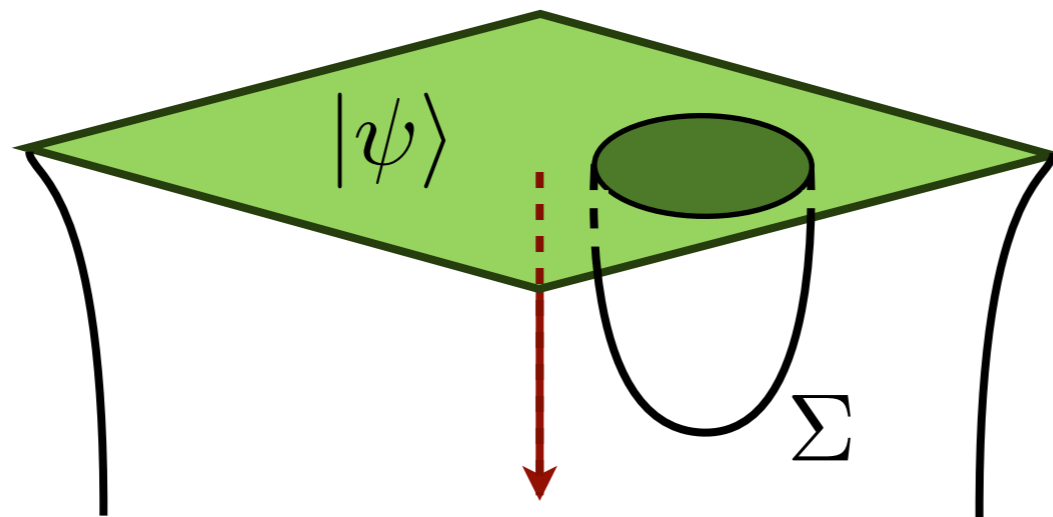
# Refining Generalized Entropy

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DAMTP, Cambridge

based on 1412.4093, with William R. Kelly (UCSB)

# Ryu-Takayanagi

$$I = \frac{1}{16\pi G} \int \sqrt{-g} d^D x R$$



Extremal surface

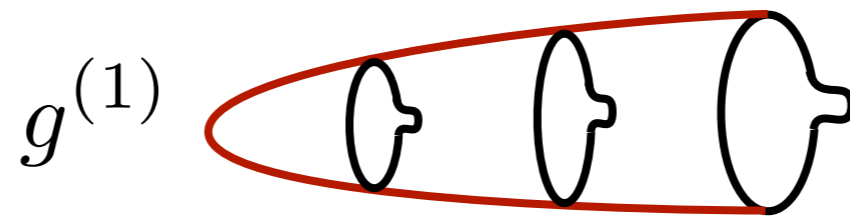
$$S = \frac{(\text{Area}(\Sigma))_{D-2}}{4G}$$

# Comments

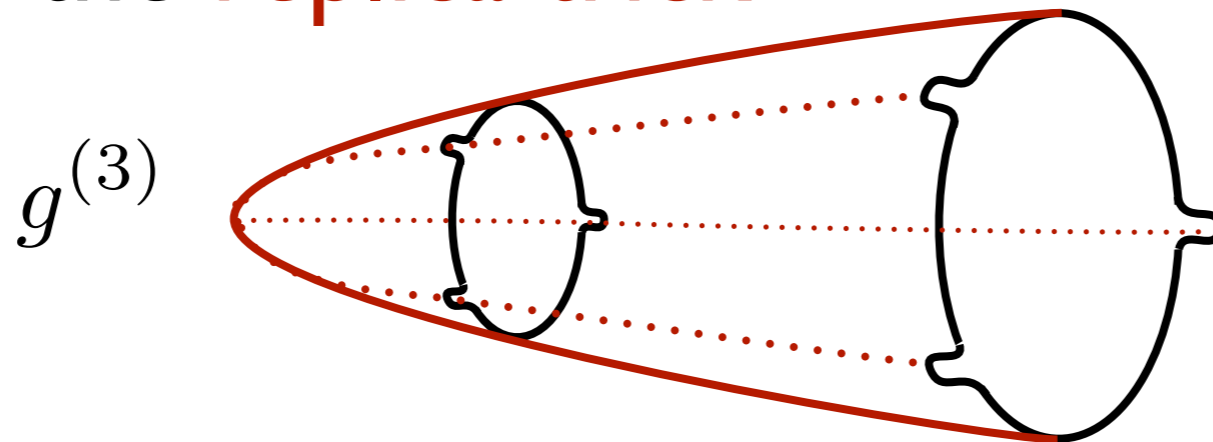
- Generalises Bekenstein-Hawking entropy
- Cumbersomely understood properties of entanglement become transparent,  $\delta S = 0$
- A tool to potentially reconstruct space holographically
- We have a first-principles derivation of it:  
Lewkowycz-Maldacena

# Outline of derivation

- Euclidean dual of thermal-like part function



- And the **replica trick**

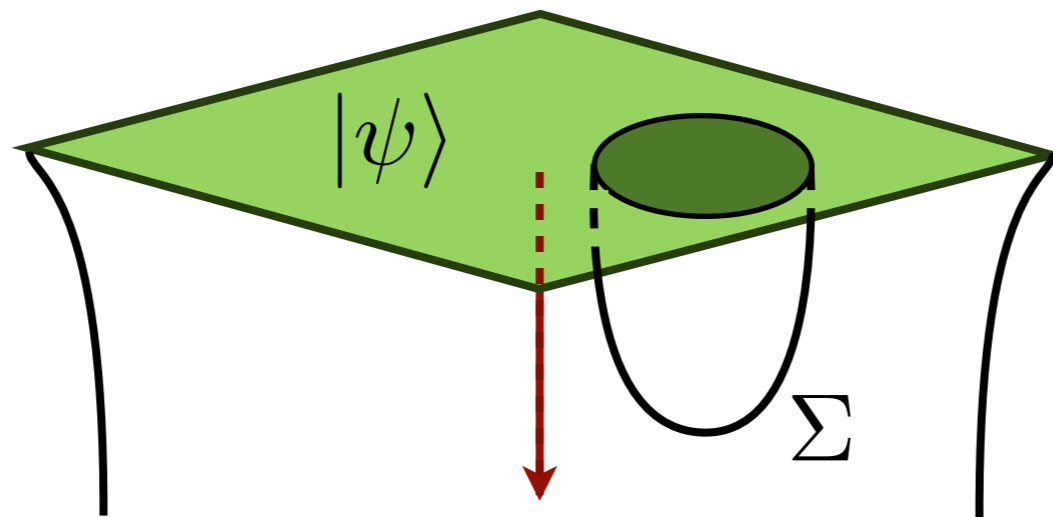


- Analytically continuing metrics  $R_{\mu\nu}[g^{(n)}] = \text{finite}$
- Expanding around  $n \sim 1$

# More general proposal

Hung, Myers, Smolkin;  
de Boer, Kulaxizi, Parnachev

$$I = \frac{1}{16\pi G} \int \sqrt{-g} d^D x (R + \alpha' \text{Riem}^2 + \dots)$$



Extremises entropy?

$$S \sim \int_{\Sigma} \sqrt{\gamma} d^{D-2} \sigma \left( \frac{\partial \mathcal{L}}{\partial \text{Riem}} + \frac{\partial^2 \mathcal{L}}{\partial \text{Riem}^2} K^2 \right)$$

Dong; JC '13  
see also Miao, Guo '14

# What this talk is about

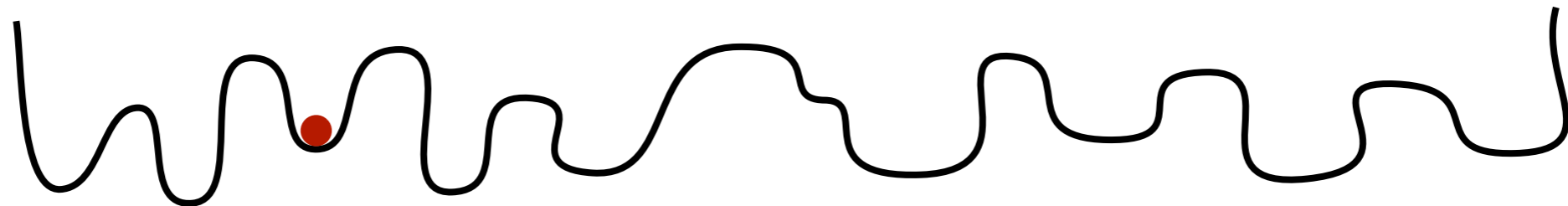
- Reviewing the replica trick in entanglement entropy calculations (and holographic dual)
- Discussing the structure of  $g^{(n)}$  (analytic continuation of metrics) around  $n \sim 1$

# Conclusions

- $\Sigma$  is determined by  $\delta S = 0$  in GR and away (explicitly in one case)
- Ryu-Takayanagi follows independently from the assumption of replica symmetry
- **Disclaimer:** We do not conclude that replica symmetry is broken (but it may be)

# Why?

- Higher derivatives: (a) Appear in the low energy limit of string theory; (b) Give a more general understanding
- Replica symmetry breaking: (a) Strengthens the derivation; (b) Hints to new possibilities
- In CMT, RSB explains glassy physics. Glasses are frustrated, with landscapes of vacua...





**Lewkowycz-Maldacena**

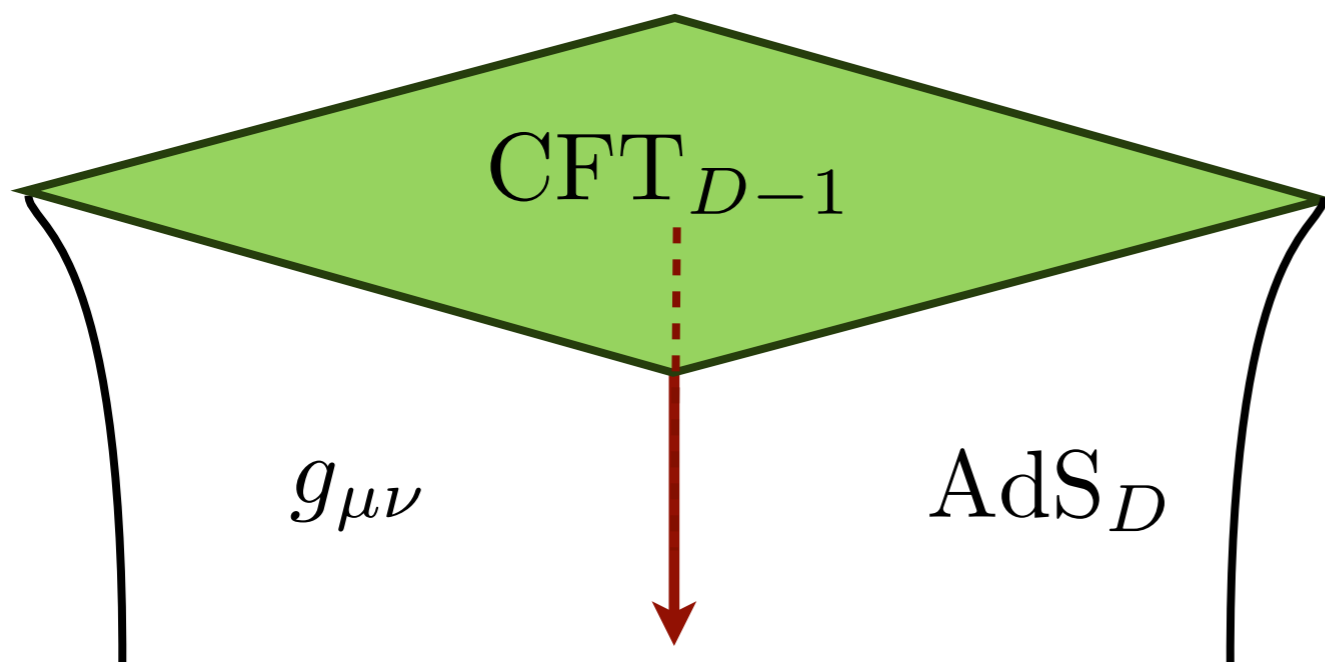
# Ingredients of Generalized Entropy

- Replica trick

$$\text{tr}(\rho \log \rho) = \lim_{n \rightarrow 1} \frac{\log(\text{tr} \rho^n)}{n - 1}$$

Expansion  
around  $n \sim 1$

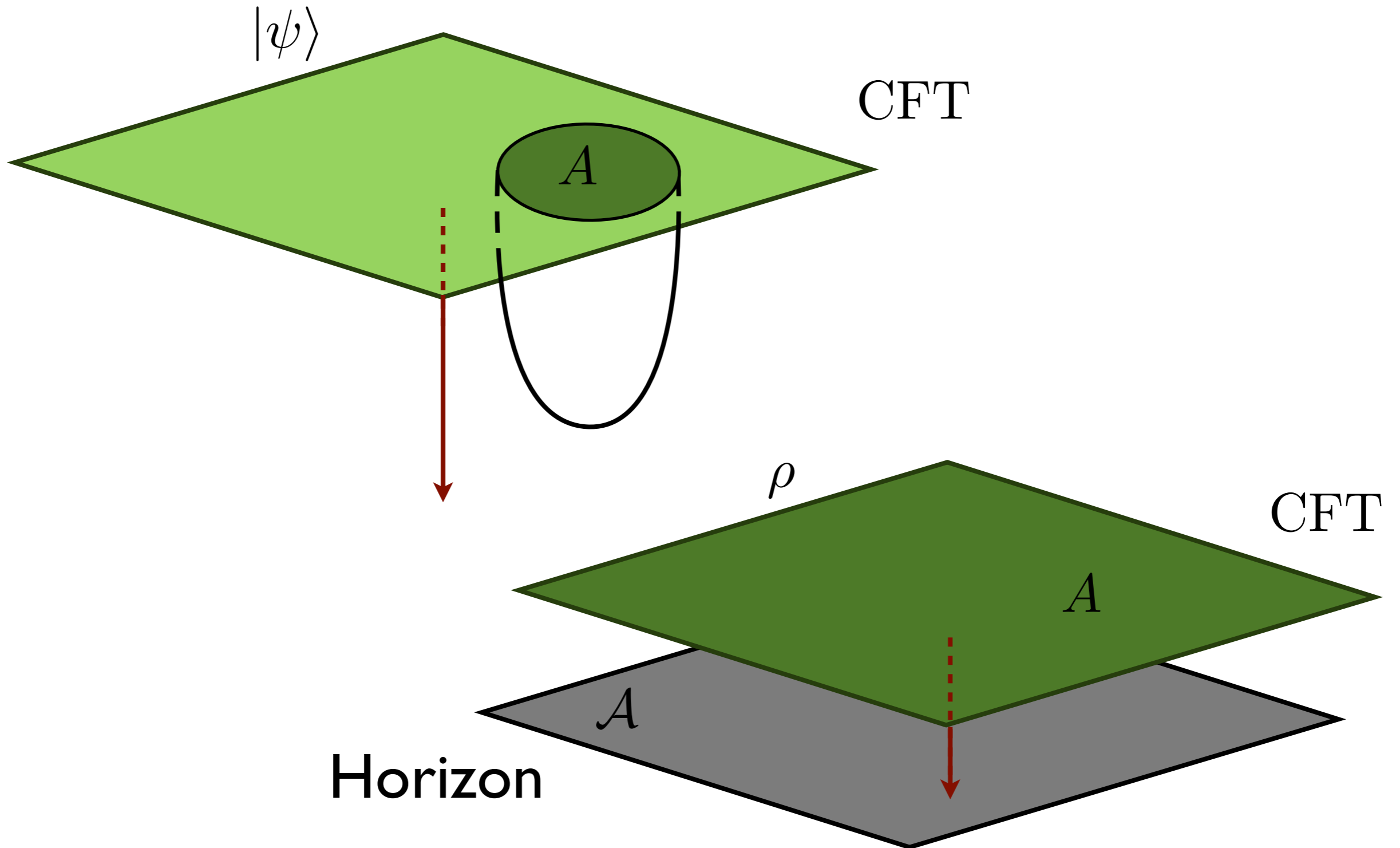
- Saddle point approximation to holography



$$Z_{\text{CFT}_{D-1}} \approx e^{-I_E[g_{\mu\nu}]}$$

# Setup

Casini, Huerta, Myers

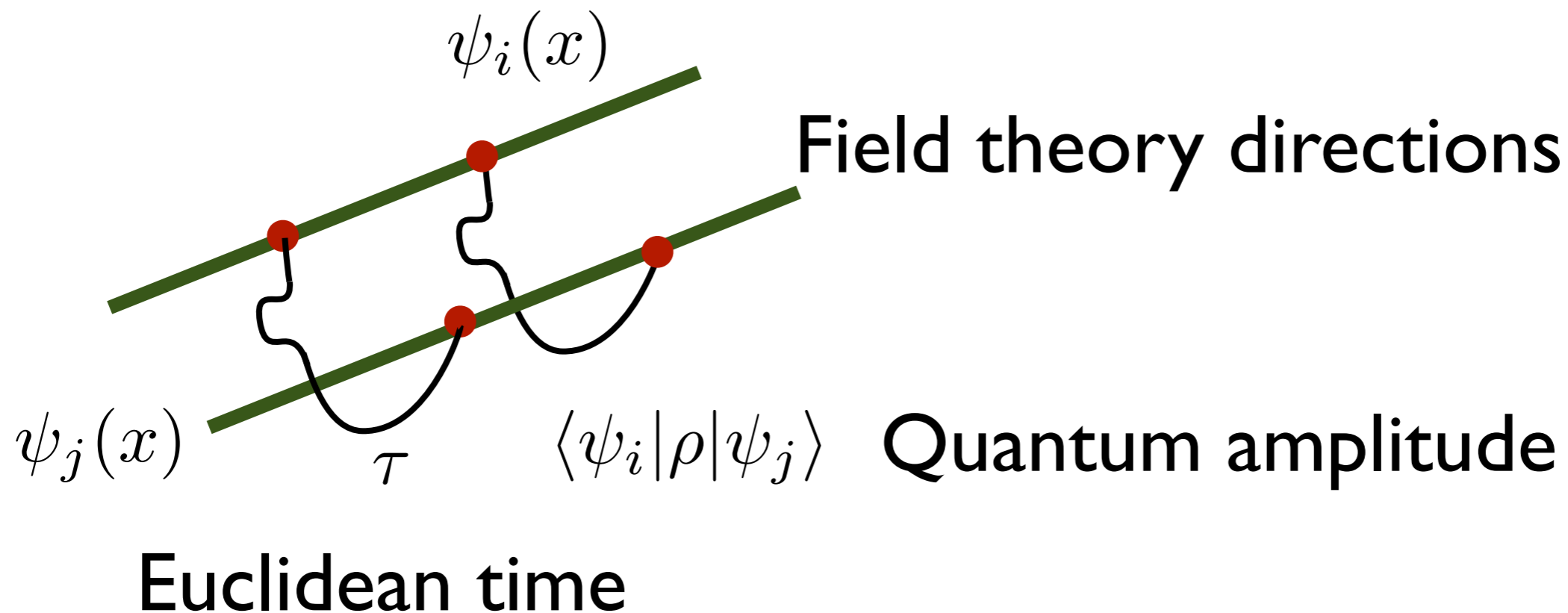


# States considered

- Generated by Euclidean path integrals

$$\rho = \mathcal{P} e^{-\int H(\tau) d\tau}$$

$$\rho_T = \sum_i |i\rangle e^{-E_i \beta} \langle i|$$

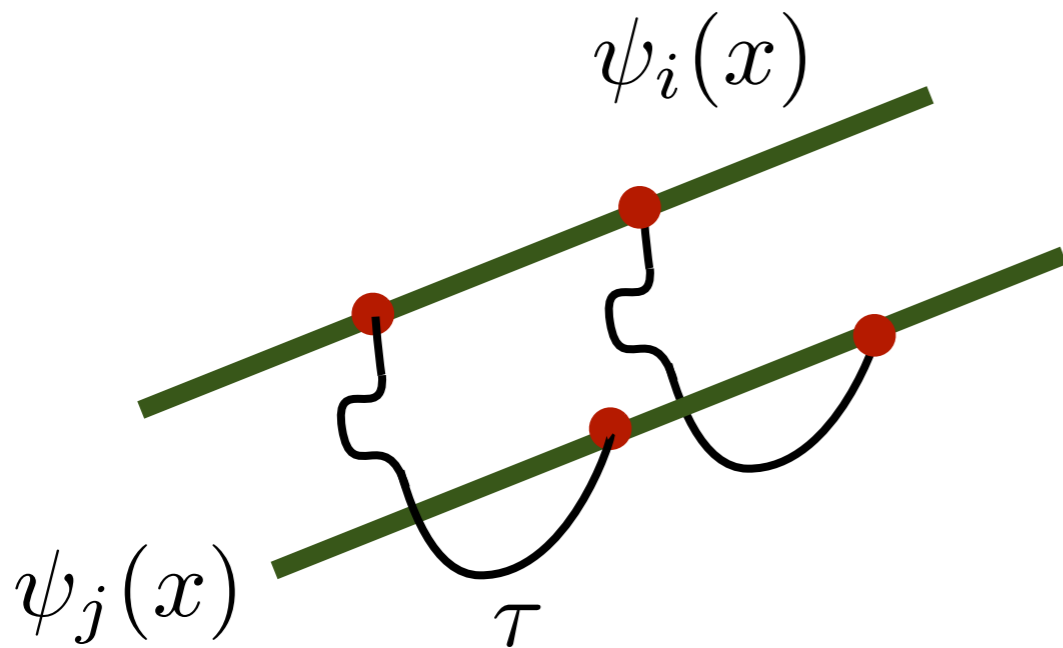


# States considered

- Generated by Euclidean path integrals

$$\rho = \mathcal{P} e^{-\int H(\tau) d\tau}$$

$$\rho_T = \sum_i |i\rangle e^{-E_i \beta} \langle i|$$



A diagram illustrating the trace of the density matrix  $\rho$ . It shows a single green line representing the boundary of a Euclidean time interval. Two black wavy loops are attached to the line, each starting and ending at a red dot on the line. The loops represent paths that start and end at the same point, forming a closed loop. Below the diagram is the equation:

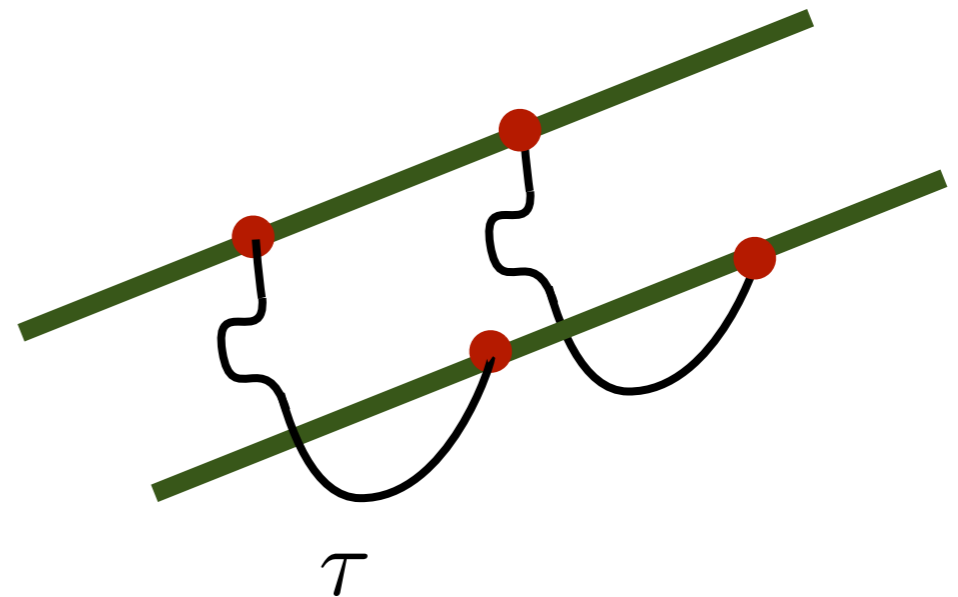
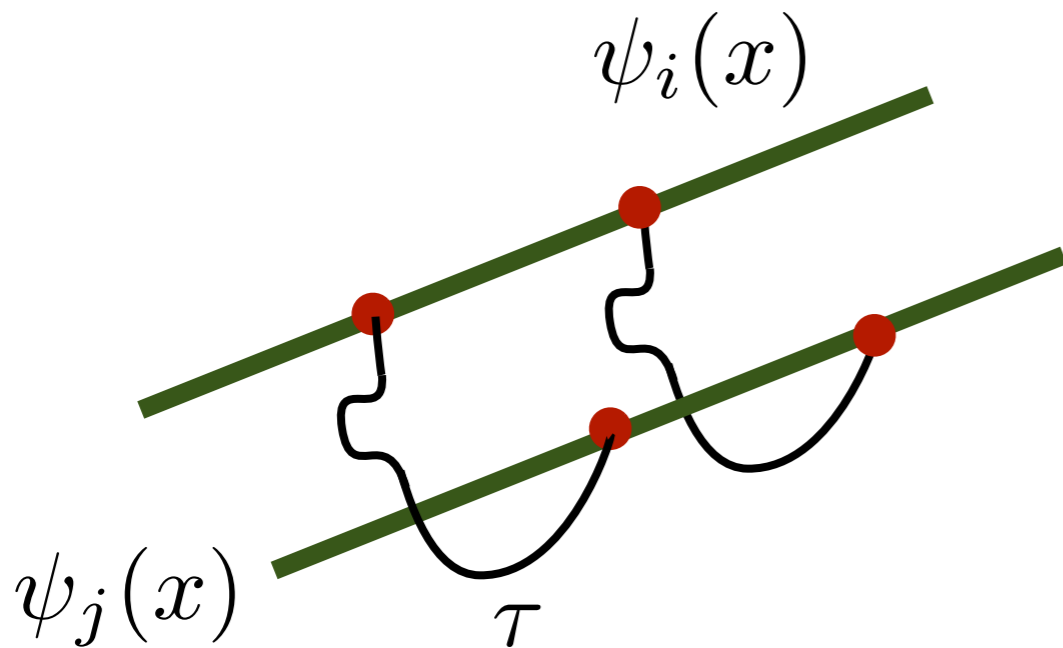
$$\sum_i \langle \psi_i | \rho | \psi_i \rangle = \text{tr} \rho$$

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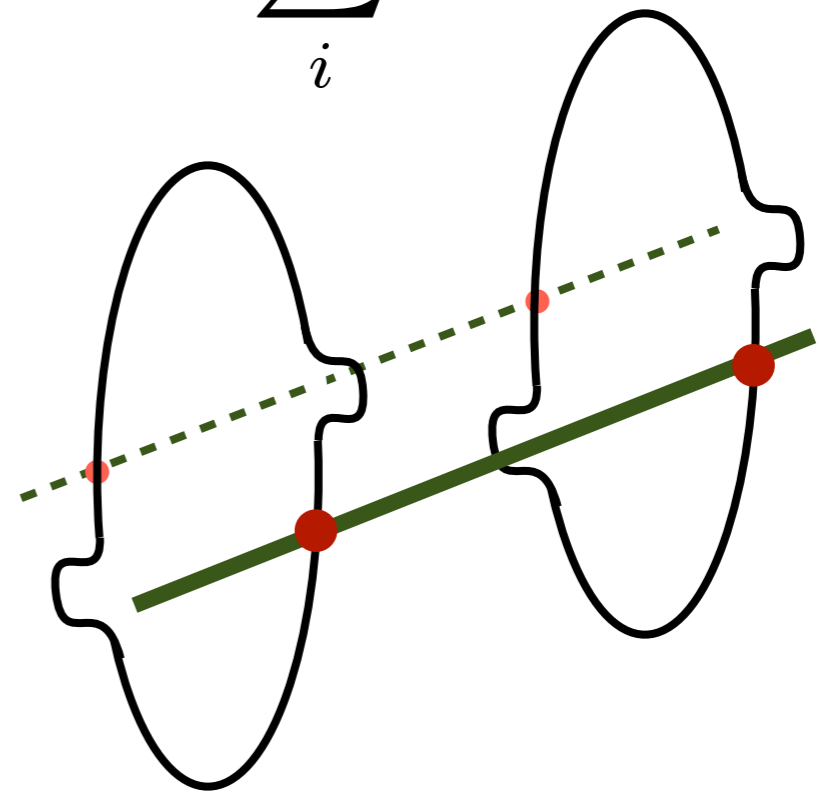
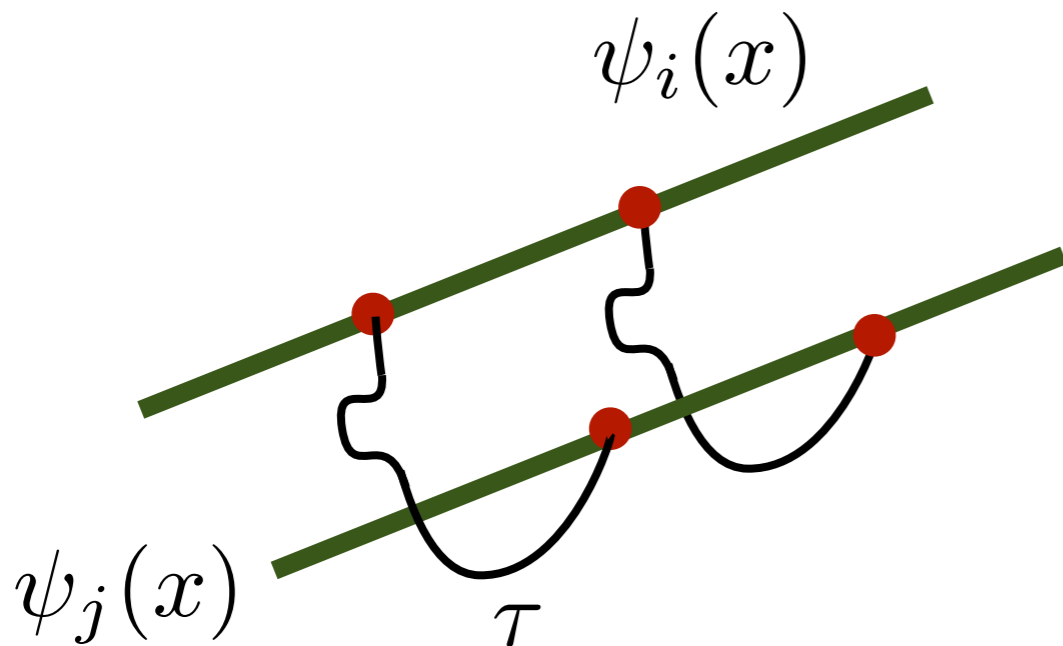
$$\text{tr} \rho^2 = \sum_i \sum_j \langle \psi_i | \rho | \psi_j \rangle \langle \psi_j | \rho | \psi_i \rangle$$

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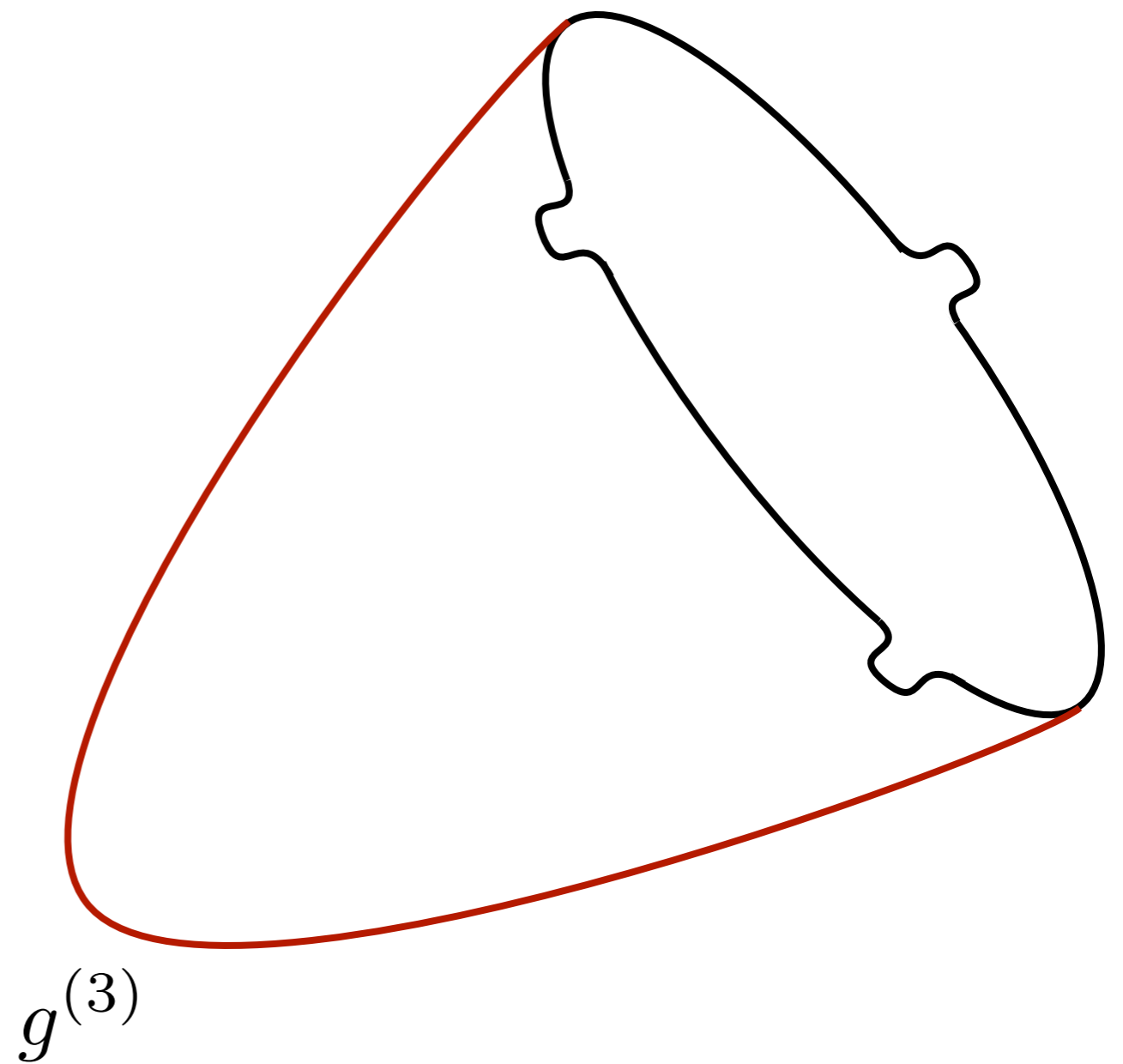
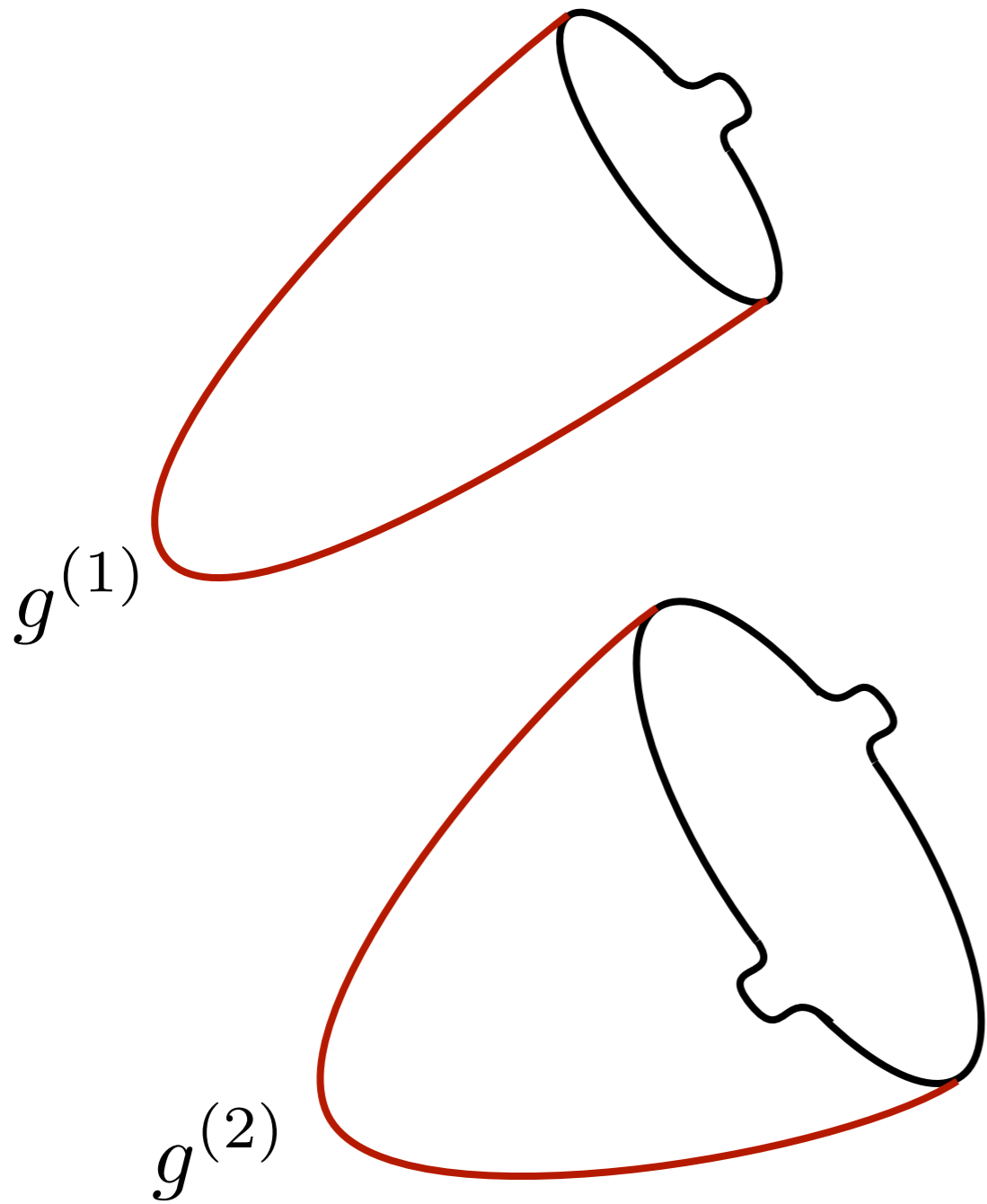
$$\rho = \mathcal{P} e^{-\int H(\tau) d\tau}$$

$$\rho_T = \sum_i |i\rangle e^{-E_i \beta} \langle i|$$



$$\text{tr} \rho^2 = \sum_i \sum_j \langle \psi_i | \rho | \psi_j \rangle \langle \psi_j | \rho | \psi_i \rangle$$

# Gravity dual



$$\frac{\text{tr}(\rho^n)}{(\text{tr}\rho)^n} \approx \frac{e^{-I[g^{(n)}]}}{e^{-nI[g^{(1)}]}}$$

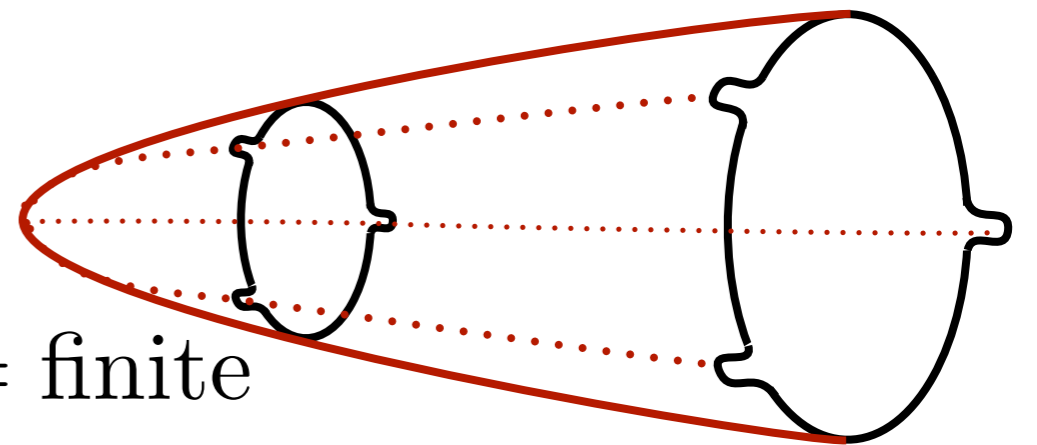


# Recap

- Replica trick  $S = \lim_{n \rightarrow 1} S_n = - \partial_n \log \frac{\text{tr}(\rho^n)}{(\text{tr} \rho)^n} \Big|_{n=1}$

- Holography

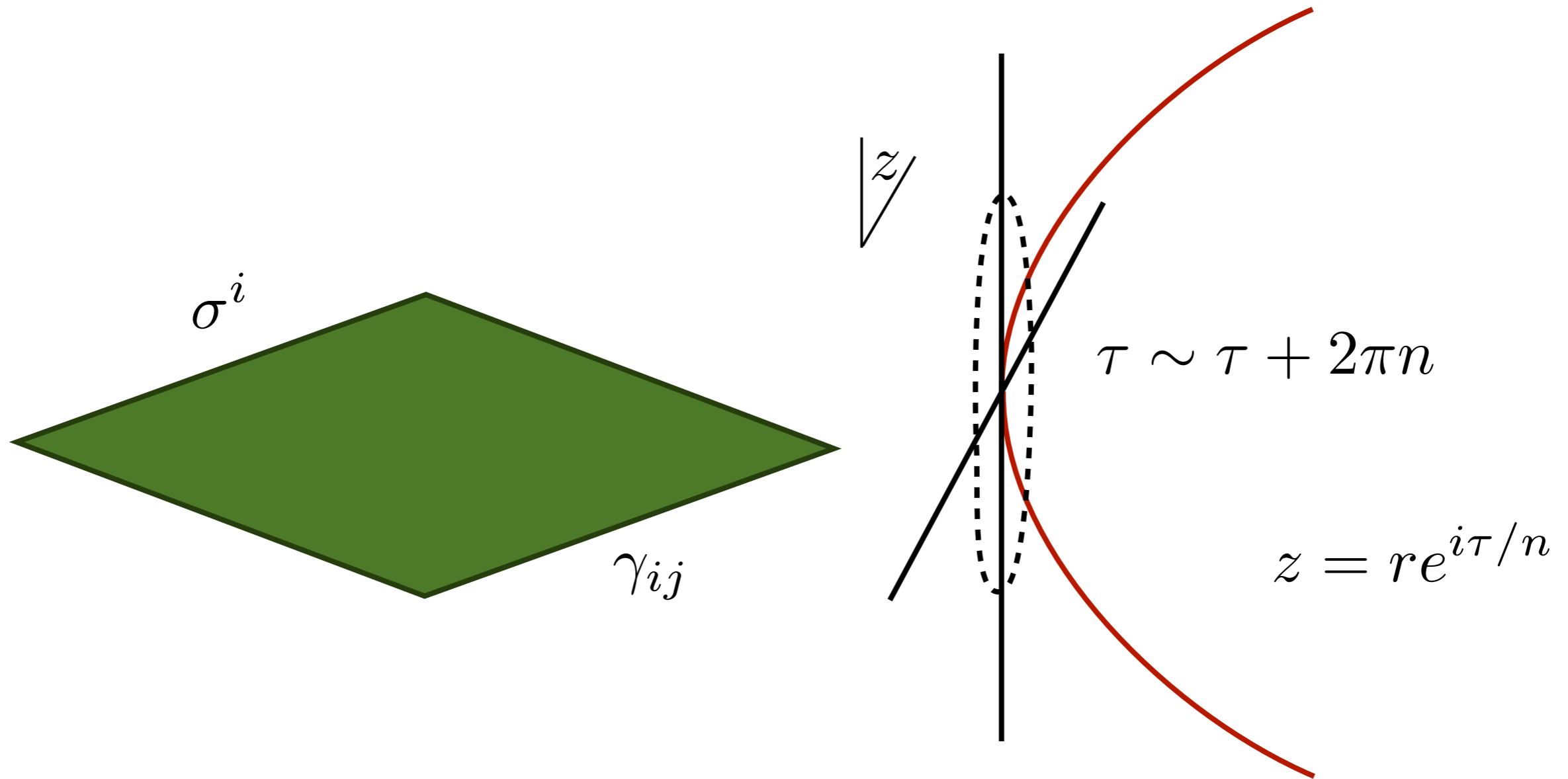
$$R_{\mu\nu} [g^{(n)}] = \text{finite}$$



- The calculation ‘localises at the tip’

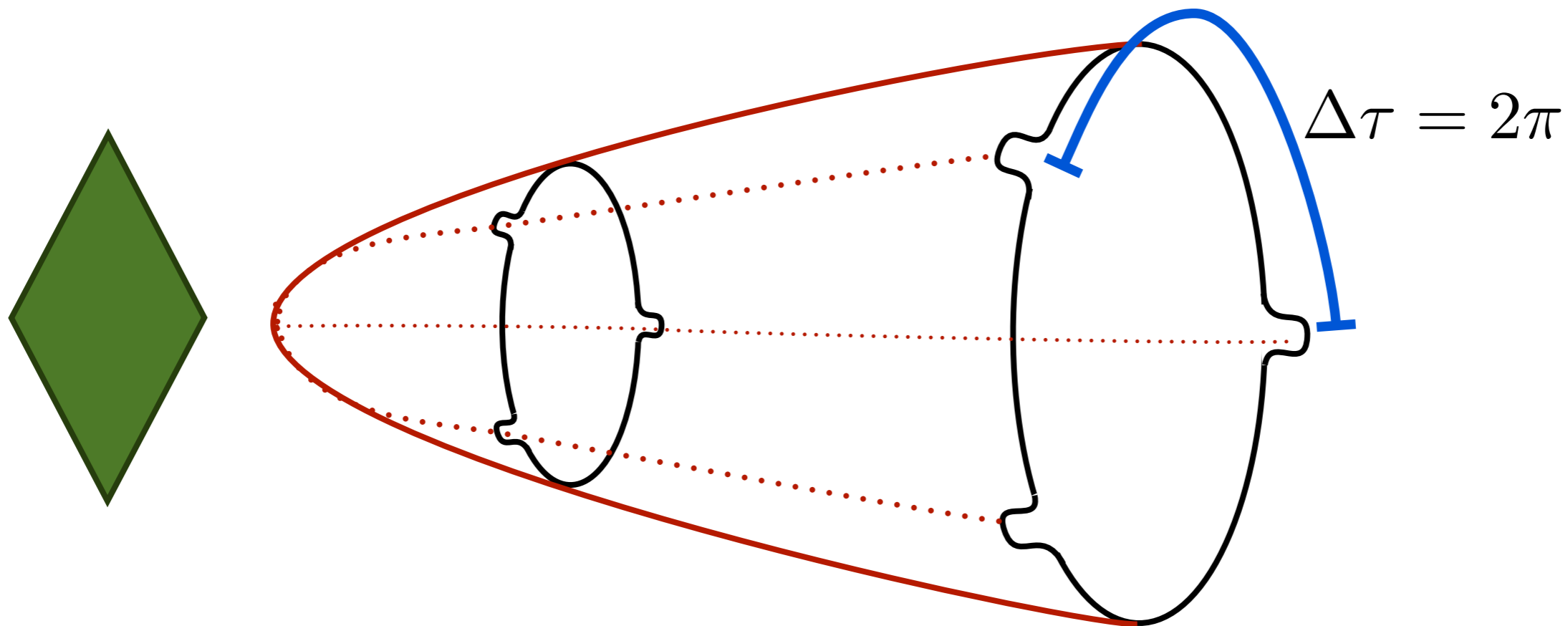
$$\mathcal{Z}_n$$

# Zooming into the tip



$$g^{(n)} = dzd\bar{z} + \gamma_{ij}d\sigma^i d\sigma^j + \dots$$

# $\mathbb{Z}_n$ symmetry



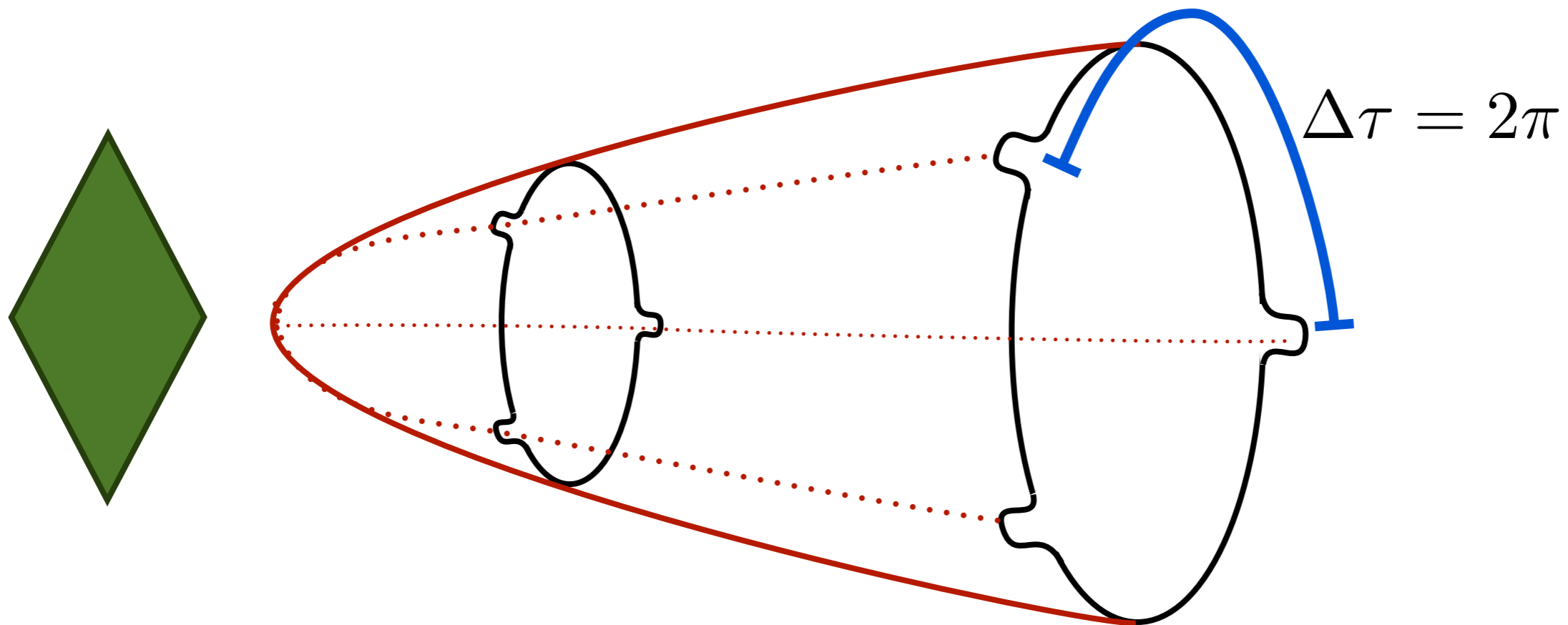
$$\tau \rightarrow \tau + 2\pi s$$

$$z = r e^{i\tau/n} \rightarrow z e^{i2\pi s/n}$$

$$z^n$$

$$z\bar{z}$$

# $\mathbb{Z}_n$ symmetry



$$g^{(n)} = dzd\bar{z} + \gamma_{ij}d\sigma^i d\sigma^j + O(z^n) + O(z\bar{z})$$

$n \sim 1$

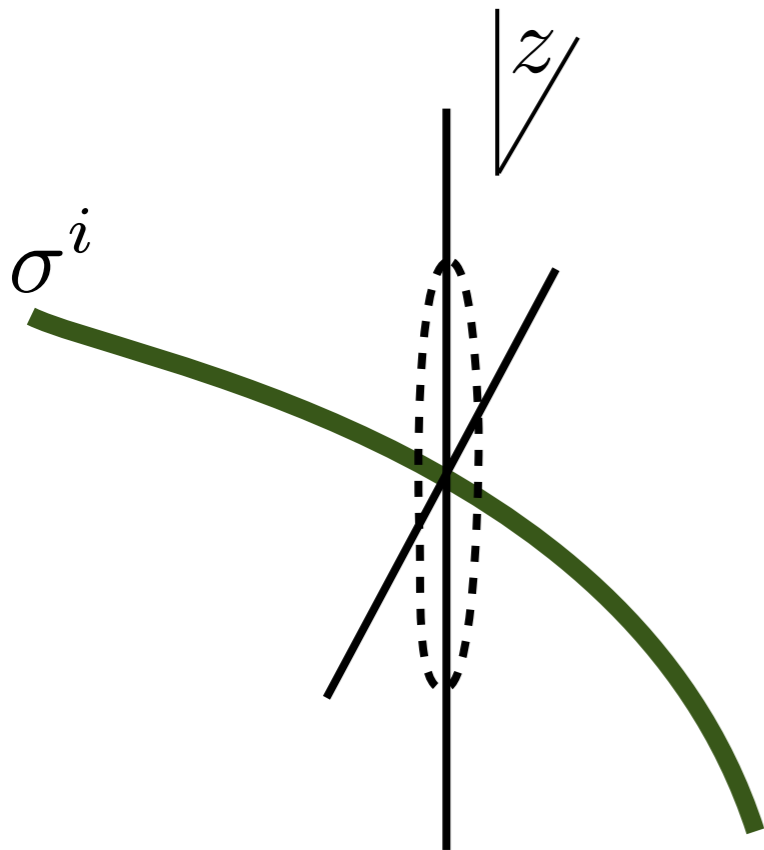
$\downarrow$   
 $O(z)$

$\downarrow$   
 $O(z^2)$

$$z^n$$

$$z\bar{z}$$

# Interlude: Geometry of surfaces



$$g = dzd\bar{z} + \gamma_{ij}d\sigma^i d\sigma^j + O(z) + O(z^2)$$

$1/\ell$

$1/\ell^2$

$R_{\mu\nu\rho\sigma}$

**Extrinsic curvature**

$$K_{ijz} z d\sigma^i d\sigma^j$$

**Torsion of curve**

$$A_{iz\bar{z}}(\bar{z}dz - zd\bar{z})d\sigma^i$$

**Gauge Christoffels**

$$\Gamma_{zz}^z z d\bar{z} dz$$

# Properties of $Z_n$ metrics

$$g^{(n)} = dzd\bar{z} + (\gamma_{ij} + 2K_{ijz}z^n) d\sigma^i d\sigma^j + \dots$$

- Explicitly regular at finite  $n$  ✓
- At finite  $n > 1$ ,  $z = 0$  has no extrinsic curvature. At  $n = 1$ ,  $K_{ijz}$  ✓
- Think about it in a double expansion

$$z \sim 0 \quad n \sim 1$$

# Properties of $Z_n$ metrics

$$g^{(n)} = dzd\bar{z} + (\gamma_{ij} + 2K_{ijz}z^n) d\sigma^i d\sigma^j + \dots$$

$$\text{Riem} \sim \partial_z^2 g^{(n)}$$

$$R_{zz} = -(n-1) \frac{\gamma^{ij} K_{ijz}}{z} z^{n-1} + \dots$$

Singular stress tensor for  $n \sim 1$  unless  $\gamma^{ij} K_{ijz} = 0$

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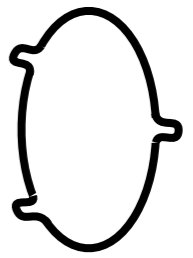


# Caveats

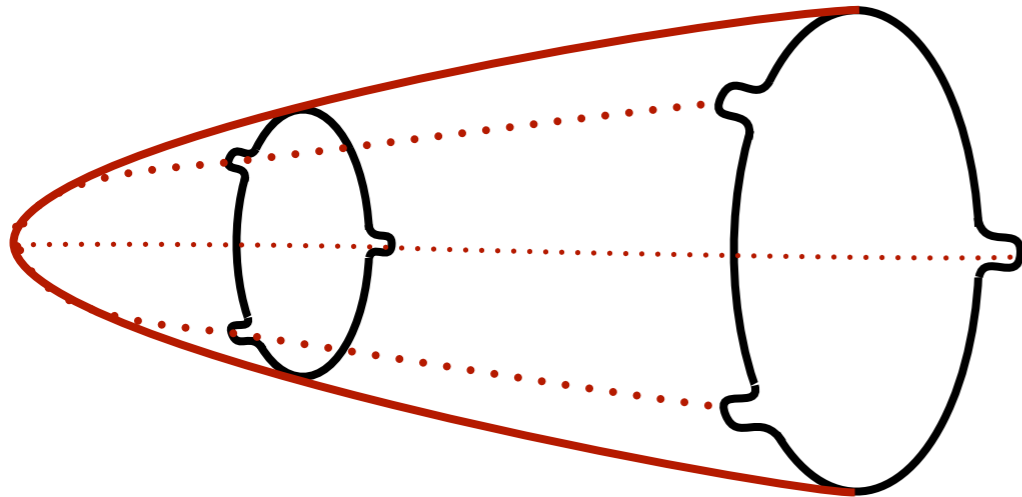
- These metrics are not single-valued for non-integer  $n$  (just as  $z^n$  is not).
- While non-single valuedness comes as a branch cut, in the Riemann tensor we see issues at just one point ( $z = 0$ ).
- This is related to the meaning of  $\mathbb{Z}_n$
- But we can calculate...

# Refinements

# RSB: discussion

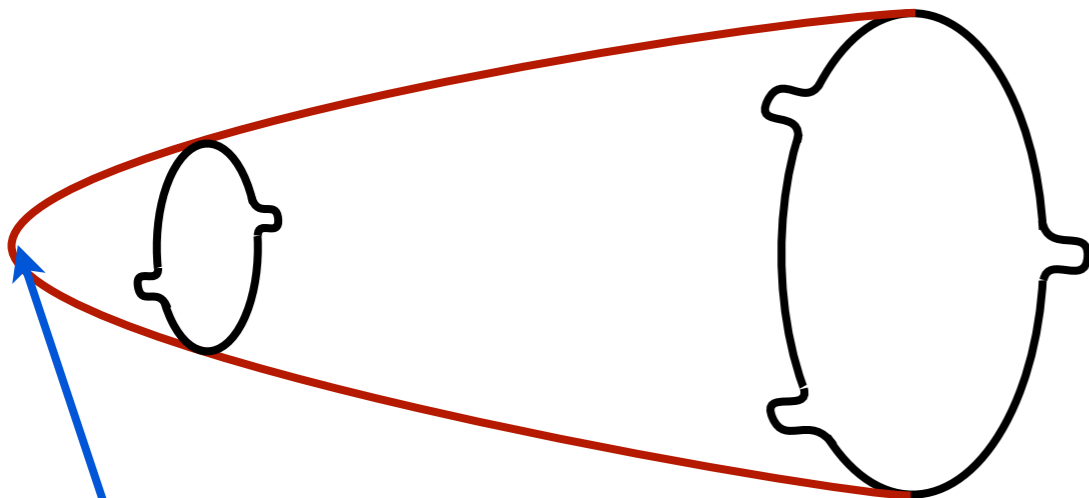
- By definition, the boundary is replica symmetric 
- The bulk could be not,  ~~$\mathbb{Z}_n$~~ , (SSB)
- The tip of the cigar can perhaps be defined as the fixed point of a residual symmetry
- However, recall we want to work to  $O(n - 1)$

# RSB terms



$$z^n$$
$$z \rightarrow z e^{i2\pi \frac{s}{n}}$$

$$\mathbb{Z}_n$$



$$z^{1+m(n-1)}$$
$$z \rightarrow z e^{i2\pi \frac{s}{1+m(n-1)}}$$

$$\mathbb{Z}_{1+m(n-1)}$$

Fixed point of residual symmetry

# RSB terms

$$K_{ijz} z^n \rightarrow \sum_{m>0} K_{ijz}^{(m)} z z^{m(n-1)}$$

- we forbid contributions  $\propto z d\sigma^i d\sigma^j$ , that do not preserve any ‘subsymmetry’ ✓
- Explicitly regular at finite  $n$  ✓
- At finite  $n > 1$ ,  $z = 0$  has no extrinsic curvature. At  $n = 1$ ,  $K_{ijz} = \sum_{m>0} K_{ijz}^{(m)}$

Geometrically located, even when RSB

# RSB terms

$$K_{ijz} z^n \rightarrow \sum_{m>0} K_{ijz}^{(m)} z z^{m(n-1)}$$

$$R_{zz} \sim \frac{n-1}{z} \sum_{m>0} m \gamma^{ij} K_{ijz}^{(m)} z^{m(n-1)}$$

$$\gamma^{ij} K_{ijz}^{(m)} = 0 \quad \longrightarrow \quad \gamma^{ij} \sum_{m>0} K_{ijz}^{(m)} = 0$$

# RSB terms

$$K_{ijz} z^n \rightarrow \sum_{m>0} K_{ijz}^{(m)} z z^{m(n-1)}$$

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$$\gamma^{ij} K_{ijz}^{(m)} = 0 \quad \longrightarrow \quad \gamma^{ij} \sum_{m>0} K_{ijz}^{(m)} = 0$$

# Further refinements

- Essential for higher derivative gravity
- Crucial to allow for terms that are pure gauge at  $n \sim 1$ , e.g.

$$g^{(n)} = dz d\bar{z} + \sum_m \Gamma_{\bar{z}\bar{z}z}^{(m)} z (z\bar{z})^{m(n-1)} dz d\bar{z} + \dots \quad \checkmark$$

- All components  $R_{\mu\nu} - \alpha' H_{\mu\nu} = \text{finite}$
- RSB allowed as before

$$\delta S = 0$$



# Summary

- Analytically continue  $g^{(n)}$  in  $n$ , expand around  $n \sim 1$



- Allow all terms in  $g^{(n)}$  consistent with hypotheses (**RSB**, pure gauge at  $n = 1$ )
- RT and generalisations follow:  $\delta S = 0$

